It is well known that with each Lax operator L

$$L\psi(x,\lambda) \equiv i\frac{d\psi}{dx} + (q(x,t) - \lambda J)\psi(x,t,\lambda) = 0$$
(1)

related to the simple Lie algebra \mathfrak{g} one can relate a hierarchy of integrable nonlinear evolution equations (NLEE)

$$iad {}_{J}^{-1} \frac{\partial q}{\partial t} + f(\Lambda)q(x,t) = 0$$
 (2)

solvable through the inverse scattering method (ISM), see the review paper [1]. Here J is a constant element of the Cartan subalgebra \mathfrak{h} of \mathfrak{g} , the potential q(x, t) takes values in $\mathfrak{g}/\mathfrak{h}$, $\Lambda = (\Lambda_+ + \Lambda_-)/2$ is the generating operator:

$$\Lambda_{\pm}X = \operatorname{ad}_{J}^{-1} \left(\operatorname{i}\frac{\mathrm{d}X}{\mathrm{d}x} + P_{0}([q(x), X(x)] + \operatorname{i}\left[q(x), \int_{\pm\infty}^{x} \mathrm{d}y[q(y), X(y)]\right] \right)$$
(3)

and $P_0 \cdot = \operatorname{ad}_J^{-1} \operatorname{ad}_J \cdot$. A most effective method to solve the inverse scattering problem for L is based on a local Riemann-Hilbert problem [2]

$$\xi^{+}(x,\lambda) = \xi^{-}(x,\lambda)G(x,\lambda), \qquad \lambda \in \Gamma$$
(4)

on the contour Γ in the complex λ -plane; in fact Γ defines the continuous spectrum of L and $G(x, \lambda)$ defines the spectral data of L. We show that to each RHP one can relate generalized exponentials and a recursion operator Λ which play fundamental role in the theory of the NLEE.

We review a class of RHP with additional symmetry conditions which can be used to solve NLEE with additional reductions. This method can be applied to a wide class of Lax operators with vanishing and non-vanishing boundary conditions. For the special case of \mathbb{Z}_h reduction group [3] where h is the Coxeter number of \mathfrak{g} we derive the action-angle variables of the corresponding NLEE, which include the two-dimensional Toda field theories.

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