

Kepler Problem and Formally Real Jordan Algebras I

Kepler problem and Lorentz transformations

Based on [G. Meng, J. Math. Phys. **53**, 052901(2012)]

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God is a mathematician of a very high order — P. Dirac

What is Kepler Problem?

- It is a mathematical model for the simplest solar system.
- I. Newton introduced and solved it in 1678, and that leads to a good explanation for Kepler's three laws of planetary motion.
- It is also a mathematical model for the simplest atom (i.e. the hydrogen atom).
- E. Schrödinger introduced and solved it (at the quantum level) in 1925, and that leads to a good explanation for the spectral lines of the hydrogen gas and Mendeleev's periodic table for elements as well.
- It is a classical example of combining beauty, simplicity and truth all in one.

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The
core of
beauty is
simplicity.

- Paulo Coelho

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Simplicity is
the seal of
truth!

Herman Boerhaave

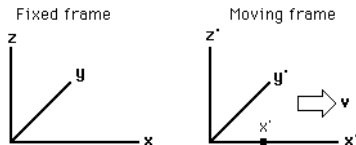
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What are Lorentz Transformations?

- They are the linear transformations of the form

$$t = \gamma(t' + \frac{vx'}{c^2}), \quad x = \gamma(x' + vt'), \quad y = y', \quad z = z'$$

for the time and (rectangular) space coordinates in two inertial frames:



Here c is the speed of light and $\gamma = \frac{1}{\sqrt{1 - \frac{v^2}{c^2}}}$.

- They are the correction to the Galileo transformations:

$$t = t', \quad x = x' + vt', \quad y = y', \quad z = z'.$$

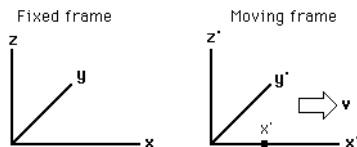
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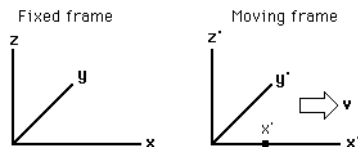
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- They (not Galileo transformations) leave invariant of the form of Maxwell equations for electromagnetism.
- The attempt to understand their geometric/physical meaning led to the relativity revolution in the early 20th century, including, among other things, the relativistic correction to Newtonian mechanics as well as various relativistic corrections to the Universal Gravitation Law, with Einstein's General Theory of Relativity being the favorite one.
- First published by Joseph Larmor in 1897 and independently again by Hendrik Antoon Lorentz in 1899.
- In mathematics, any linear transformation $T: \mathbb{R}^4 \rightarrow \mathbb{R}^4$ that preserves the Lorentz inner product:

$$(x_0, \mathbf{x}) \cdot (y_0, \mathbf{y}) = x_0 y_0 - \mathbf{x} \cdot \mathbf{y}$$

is called a Lorentz transformation. Then Lorentz transformations form a group, i.e., the Lie group $O(1, 3)$. For simplicity, we shall write x for (x_0, \mathbf{x}) , x^2 for $x \cdot x$. So $x^2 = x_0^2 - \mathbf{x} \cdot \mathbf{x}$.

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- A new description of the orbits
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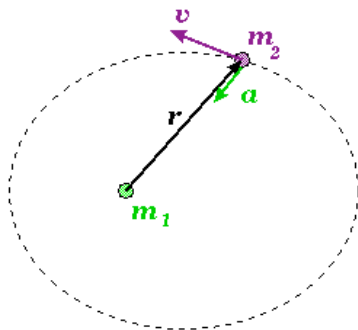
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A brief review of the Kepler problem

- Configuration space: $\mathbb{R}_*^3 := \mathbb{R}^3 \setminus \{\mathbf{0}\}$.
- Equation of Motion:

$$\mathbf{r}'' = -\frac{\mathbf{r}}{r^3}. \quad (1)$$



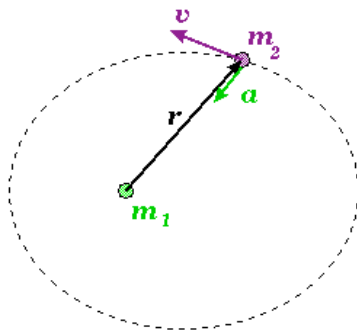
- Angular Momentum $\mathbf{L} := \mathbf{r} \times \mathbf{r}'$ is conserved:

$$\mathbf{L}' = \mathbf{r}' \times \mathbf{r}' + \mathbf{r} \times \mathbf{r}'' = \mathbf{r} \times \left(-\frac{\mathbf{r}}{r^3}\right) = \mathbf{0}.$$

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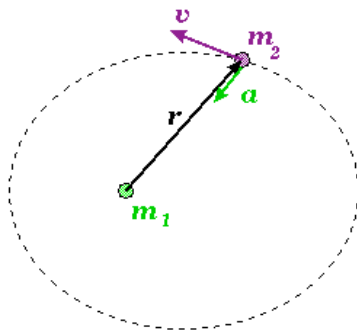
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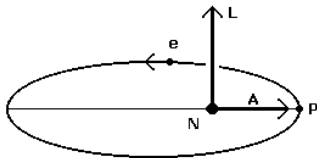
- Lenz vector $\mathbf{A} := \mathbf{L} \times \mathbf{r}' + \frac{\mathbf{r}}{r}$ is conserved:

$$\begin{aligned}\mathbf{A}' &= \mathbf{L} \times \mathbf{r}'' + \left(\frac{\mathbf{r}}{r}\right)' = -(\mathbf{r} \times \mathbf{r}') \times \frac{\mathbf{r}}{r^3} + \left(\frac{\mathbf{r}}{r}\right)' \\ &= -\frac{r^2 \mathbf{r}' - r r' \mathbf{r}}{r^3} + \left(\frac{\mathbf{r}}{r}\right)' = \mathbf{0}.\end{aligned}$$

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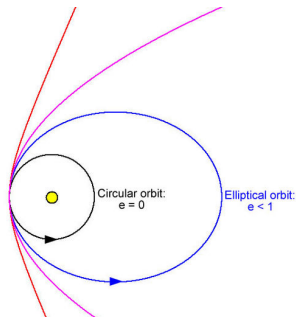
- Orbits. Since $\mathbf{L} = \mathbf{r} \times \mathbf{r}'$, $\mathbf{A} = \mathbf{L} \times \mathbf{r}' + \frac{\mathbf{r}}{r}$, we have $\mathbf{L} \cdot \mathbf{A} = 0$



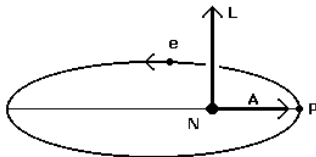
and

$$\mathbf{L} \cdot \mathbf{r} = 0, \quad r - \mathbf{A} \cdot \mathbf{r} = |\mathbf{L}|^2, \quad (2)$$

So a non-colliding orbit is a conic with eccentricity e equal to $|\mathbf{A}|$:



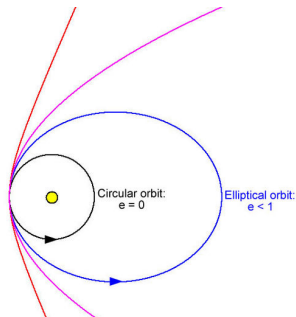
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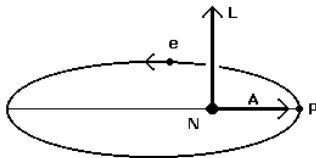
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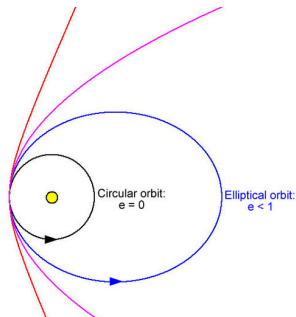
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So a non-colliding orbit is a conic with eccentricity e equal to $|\mathbf{A}|$:



- Total energy. Assume the orbit is non-colliding (i.e. $\mathbf{L} \neq \mathbf{0}$), then the total energy $E := \frac{1}{2}|\mathbf{r}'|^2 - \frac{1}{r}$ can be expressed in terms of \mathbf{L} and \mathbf{A} :

$$E = -\frac{1 - |\mathbf{A}|^2}{2|\mathbf{L}|^2}. \quad (3)$$

Proof.

$$\begin{aligned} |\mathbf{A}|^2 &= |\mathbf{L} \times \mathbf{r}'|^2 + 2\frac{\mathbf{r} \cdot (\mathbf{L} \times \mathbf{r}')}{r} + 1 \\ &= |\mathbf{L}|^2|\mathbf{r}'|^2 - 2\frac{|\mathbf{L}|^2}{r} + 1 \\ &= 2|\mathbf{L}|^2E + 1. \end{aligned}$$

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Magnetized Kepler Problems

- Magnetized Kepler problems were introduced towards the end of 1960s, by H. McIntosh and A. Cisneros, and independently by D. Zwanziger, so they are called **MICZ-Kepler problems**.
- They are the mathematical models for the hypothetical hydrogen atoms for which the nucleus carries magnetic charge as well.
- Their configuration spaces are all the same: $\mathbb{R}_*^3 := \mathbb{R}^3 \setminus \{\mathbf{0}\}$.
- For the hypothetical hydrogen atom whose nucleus carries magnetic charge μ , its equation of motion is

$$\mathbf{r}'' = -\frac{\mathbf{r}}{r^3} - \mathbf{r}' \times \mu \frac{\mathbf{r}}{r^3} + \frac{\mu^2}{r^4} \mathbf{r} \quad (4)$$

Conserved quantities are angular momentum $\mathbf{L} := \mathbf{r} \times \mathbf{r}' + \mu \frac{\mathbf{r}}{r}$ and Lenz vector $\mathbf{A} := \mathbf{L} \times \mathbf{r}' + \frac{\mathbf{r}}{r}$.

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- It is easy to see that $\mathbf{L} \cdot \mathbf{A} = \mu$, and

$$\mathbf{L} \cdot \mathbf{r} = \mu r, \quad r - \mathbf{A} \cdot \mathbf{r} = |\mathbf{L}|^2 - \mu^2. \quad (5)$$

- Eq. (5) gives an algebraic description for the orbits, from which, we deduce that there are four types of orbits: linear, elliptic, parabolic, and hyperbolic.

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A new description for the orbits

The preceding set of algebraic equations can be rewritten as

$$\mu r - \mathbf{L} \cdot \mathbf{r} = 0, \quad r - \mathbf{A} \cdot \mathbf{r} = |\mathbf{L}|^2 - \mu^2. \quad (6)$$

Assume that the orbit is non-colliding, i.e. $|\mathbf{L}|^2 - \mu^2 = |\mathbf{r} \times \mathbf{r}'|^2 > 0$. Then, we can introduce 4-D Lorentz vectors

$$l = \frac{1}{\sqrt{|\mathbf{L}|^2 - \mu^2}}(\mu, \mathbf{L}), \quad a = \frac{1}{|\mathbf{L}|^2 - \mu^2}(1, \mathbf{A}), \quad x = (r, \mathbf{r}). \quad (7)$$

so that Eq. (6) can be rewritten as

$$l \cdot x = 0, \quad a \cdot x = 1. \quad (8)$$

It is easy to see that $l^2 = -1$, $l \cdot a = 0$, $a_0 > 0$, and

$$E = -\frac{a^2}{2a_0}.$$

Remark: Eq. (8) is for $\mathbf{r} \in \mathbb{R}_*^3$, but it is also for $x \in \mathbb{R}^4$ provided that x is on the future light cone.

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A new description for the orbits

The preceding set of algebraic equations can be rewritten as

$$\mu r - \mathbf{L} \cdot \mathbf{r} = 0, \quad r - \mathbf{A} \cdot \mathbf{r} = |\mathbf{L}|^2 - \mu^2. \quad (6)$$

Assume that the orbit is non-colliding, i.e. $|\mathbf{L}|^2 - \mu^2 = |\mathbf{r} \times \mathbf{r}'|^2 > 0$. Then, we can introduce 4-D Lorentz vectors

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$$l \cdot x = 0, \quad a \cdot x = 1. \quad (8)$$

It is easy to see that $l^2 = -1$, $l \cdot a = 0$, $a_0 > 0$, and

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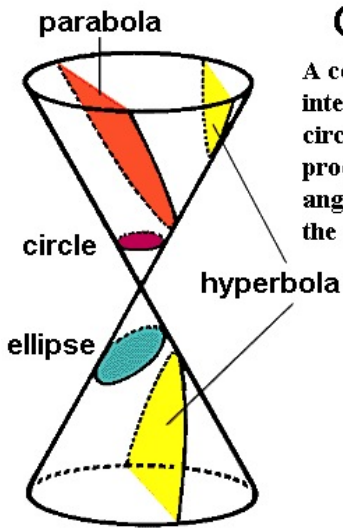
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Future light cone in the 3-D Lorentz space



Conic Sections

A conic section is formed by the intersection of a plane with a right circular cone. The "kind" of curve produced is determined by the angle at which the plane intersects the surface.

Kepler Problem and Lorentz Transformations

MICZ-Kepler orbit — a non-colliding orbit in a MICZ-Kepler problem. There are three types: elliptic, parabolic, and hyperbolic.

small Lorentz transformation — a small linear transformation T from \mathbb{R}^4 to \mathbb{R}^4 which preserves the Lorentz inner product. Here “small” means that T can be continuously deformed to the identity map on \mathbb{R}^4 .

scaling transformation — the scalar multiplication of vectors in \mathbb{R}^4 by a positive real number.

Theorem (G. Meng, J. Math. Phys. **53**, 052901(2012))

- 1) Any two oriented parabolic MICZ-Kepler orbits can be transformed from one to the other via a little Lorentz transformation.
- 2) Any two oriented elliptic MICZ-Kepler orbits can be transformed from one to the other via a little Lorentz transformation together with a scaling transformation.

Remark. 1) A second temporal dimension (i.e. x_0) appears naturally.
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